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Electron collisions with phenol: Total, integral, differential, and momentum transfer cross sections and the role of multichannel coupling effects on the elastic channel

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We report theoretical and experimental total cross sections for electron scattering by phenol (C_6H_5OH). The experimental data were obtained with an apparatus based in Madrid and the calculated cross sections with two different methodologies, the independent atom method with screening corrected additivity rule (IAM-SCAR), and the Schwinger multichannel method with pseudopotentials (SMCPP). The SMCPP method in the N_{open} -channel coupling scheme, at the static-exchange-plus-polarization approximation, is employed to calculate the scattering amplitudes at impact energies ranging from 5.0 eV to 50 eV. We discuss the multichannel coupling effects in the calculated cross sections, in particular how the number of excited states included in the open-channel space impacts upon the convergence of the elastic cross sections at higher collision energies. The IAM-SCAR approach was also used to obtain the elastic differential cross sections (DCSs) and for correcting the experimental total cross sections for the so-called forward angle scattering effect. We found a very good agreement between our SMCPP theoretical differential, integral, and momentum transfer cross sections and experimental data for benzene (a molecule differing from phenol by replacing a hydrogen atom in benzene with a hydroxyl group). Although some discrepancies were found for lower energies, the agreement between the SMCPP data and the DCSs obtained with the IAM-SCAR method improves, as expected, as the impact energy increases. We also have a good agreement among the present SMCPP calculated total cross section (which includes elastic, 32 inelastic electronic excitation processes and ionization contributions, the latter estimated with the binary-encounter-Bethe model), the IAM-SCAR total cross section, and the experimental data when the latter is corrected for the forward angle scattering effect [Fuss *et al.*, Phys. Rev. A **88**, 042702 (2013)]. © 2015 AIP Publishing LLC. [<http://dx.doi.org/10.1063/1.4913824>]

I. INTRODUCTION

One of the relevant challenges for a theoretician working on electron-molecule scattering calculations consists of defining a proper balance between the open- and closed-channel electronic space, for description of the $(N + 1)$ -electron problem to be solved. Separation of the nuclear motion from the scattering dynamics caused by the electronic cloud in a fixed-nuclei approximation is the usual strategy to simplify the problem. Even at this level of approximation, in molecular targets, the density of electronically excited states is usually so large that the convergence of multichannel scattering calculations

could be seriously hindered in many cases of practical interest. This is particularly true for molecules having its first thresholds opening up at 3–4 eV (as is the case of the DNA bases) and becomes even more critical as the energy of the incident electron increases. In such cases it is therefore necessary to judiciously truncate the electronic configuration space. At sufficiently low energies, only the elastic channel is open and the distortion of the electronic cloud is taken into account by allowing virtual excitations (closed channel space) from the ground state. This effect combined with the proper solution of the scattering problem, including the exchange interaction, is known as static-exchange-plus-polarization (SEP)

approximation. As we increase the impact energy, several discrete and continuum states (above the ionization threshold) can be excited giving rise to the important questions: (1) how many of these states must be included in a calculation in order to provide a reasonable approximation of the experimental total cross section (TCS)? (2) how do these multichannel effects affect the elastic differential cross sections (DCSs) and their related integral and momentum transfer cross sections (MTCS)? (3) how do they affect the other electronic excitation cross sections? and (4) how sophisticated must the electronic states description be in order to accurately calculate the cross sections? Recent applications¹⁻⁴ performed by theoretical groups using different methods for describing the electronic excitation of molecules by electron impact, at specific levels of multichannel coupling, have addressed one or more of these issues. They clearly indicate that there are many challenges that must be overcome to obtain reliable cross sections. On the experimental side, the energy separation of the electronic states is the main challenge. It is thus very common to work with unresolved states, particularly in polyatomic species, and pursue measurements for bands containing a few or more states related to a particular energy loss range.

In this paper, we present a theoretical and experimental effort to obtain cross sections for scattering of low and intermediate energy electrons by the phenol molecule. Phenol is an interesting molecule for technological applications⁵ and in addition it contains several characteristics that make this joint effort relevant: (1) it is a small complex organic molecule, (2) it has a high density of electronic states at relatively small impact energies, and (3) it has a permanent dipole moment, which is always a challenge for both total cross section measurements and the calculations. The present paper is outlined as follows. In Secs. II–V, different aspects of our theoretical models and the corresponding computational details are briefly described. In Sec. VI, experimental details of the apparatus in Madrid, used to obtain the present total cross sections, are described. Results and a discussion of these results are subsequently presented in Secs. VII and VIII, with a small comment on the choice of ℓ_{max} in the calculations being given in Sec. IX. Finally, in Sec. X, some concluding remarks from this investigation are summarized.

II. THE SCHWINGER MULTICHANNEL METHOD WITH PSEUDOPOTENTIALS (SMCPP)

The Schwinger multichannel method (SMC)⁶ is a variational approach that uses square integrable basis functions to obtain the scattering amplitude for electron-molecule collisions, including important effects such as exchange, polarization, and electronic multichannel coupling. In the present calculations, we use the parallel version⁷ of the SMC implementation that employs norm-conserving pseudopotentials⁸ (SMCPP) and single-excitation configuration interaction techniques for the target description.⁹ Since the method and its computational implementation were described in detail elsewhere,⁹ here we only give the working expression for the scattering amplitude

$$f(\mathbf{k}_f, \mathbf{k}_i) = -\frac{1}{2\pi} \sum_{m,n} \langle S_{\mathbf{k}_f} | V | \chi_m \rangle (d^{-1})_{mn} \langle \chi_n | V | S_{\mathbf{k}_i} \rangle, \quad (1)$$

where

$$d_{mn} = \langle \chi_m | \left[\frac{\hat{H}}{N+1} - \frac{\hat{H}P + P\hat{H}}{2} + \frac{PV + VP}{2} - VG_P^{(+)}V \right] | \chi_n \rangle. \quad (2)$$

In the expressions above, P is a projector onto N_{open} energy-allowed target electronic channels, i.e.,

$$P = \sum_{\ell=1}^{N_{open}} | \Phi_\ell \rangle \langle \Phi_\ell |, \quad (3)$$

$G_P^{(+)}$ is the free-particle Green's function projected onto the P space, V is the projectile-target interaction potential, $\mathbf{k}_i$ ($\mathbf{k}_f$) is the incoming (outgoing) projectile wave vector, and $\hat{H} = E - H$ is the total energy (ground state energy plus kinetic energy of the incoming electron) minus the Hamiltonian of the $(N+1)$ electrons under the field of the fixed nuclei. The latter is given by $H = H_0 + V$, where H_0 describes the non-interacting electron-molecule system and $S_{\mathbf{k}}$ is a solution of H_0 , namely, the product of a plane wave (projectile) and a target state Φ_ℓ (obtained within the scope of the single excitation configuration-interaction description). For the expansion of the variational scattering wave function, the method employs trial basis sets comprising $(N+1)$ -particle configuration state functions (CSFs), denoted by χ_m , that are built from spin-adapted, anti-symmetrized products of target electronic states and projectile scattering orbitals. The open electronic collision channels are included in the P space, and the dynamical response of the target electrons to the projectile field (correlation-polarization effects) is accounted for through virtual excitations of the target. In this case, the CSFs are given by

$$| \chi_m \rangle = \mathcal{A}_{N+1} | \Phi_i(1, \dots, N) \rangle \otimes | \varphi_j(N+1) \rangle, \quad (4)$$

where for $i > 0$, $| \Phi_i \rangle \equiv {}^{(2S+1)}(h_i \rightarrow p_i)$ is a singly excited state obtained by promoting one electron from a hole orbital (h_i) of the ground state $\Phi_0(1, \dots, N)$ to a particle orbital (p_i), with either singlet ($S=0$) or triplet ($S=1$) spin coupling, though only $(N+1)$ -electron configurations with total spin $S=1/2$ (doublets) are actually taken into account. If we have N_{open} states in Eq. (3), this level of calculation is called an N_{open} -channel coupling scheme at the static-exchange-plus-polarization (acronym is N_{open} ch-sep) approximation. In order to transform the scattering amplitude from the body-fixed frame (the reference frame best suited for carrying out the calculations) to the laboratory-fixed frame (the reference frame where the z -axis is aligned with the direction of the incident wave vector, i.e., $\mathbf{k}_i = k_i \hat{\mathbf{z}}$), we expand $\mathbf{k}_f$ in terms of partial waves¹⁰

$$f(\mathbf{k}_f, \mathbf{k}_i) \equiv \langle \mathbf{k}_f | f | \mathbf{k}_i \rangle = \sum_{\ell=0}^{\ell_{max}} \sum_{m=-\ell}^{\ell} \langle \mathbf{k}_f | \ell m \rangle f(\ell m, \mathbf{k}_i), \quad (5)$$

where $\langle \mathbf{k}_f | \ell m \rangle$ is a spherical harmonic that can be easily converted from the body- to the laboratory-frame and $f(\ell m, \mathbf{k}_i) = \langle \ell m | f | \mathbf{k}_i \rangle$ can be understood as the scattering amplitude of an electron entering the interaction region in a plane-wave $| \mathbf{k}_i \rangle$ and leaving it in a partial wave $| \ell m \rangle$. Although not shown

here, all SMCPP differential cross sections in this paper, over the entire energy range (5-50 eV), are numerically converged with $\ell_{\max} = 10$ (except for 50 eV that demands $\ell_{\max} = 13$) if combined with a quadrature point distribution, using a 26 Gauss-Legendre scheme for $0 \leq \theta_i \leq \pi$ and 52 points for $0 \leq \phi_i \leq 2\pi$, to describe $\mathbf{k}_i(\theta_i, \phi_i)$ in spherical coordinates. Although we have contributions from high partial waves in the scattering orbitals due to the multi-center expansion, we only employ Cartesian Gaussians (CG) of s , p , and d types on the oxygen atom and on each center for the carbon atoms. On the hydrogen atoms we have only CG functions of s and p types. As discussed in our previous applications,⁴ this makes the description of the high partial waves more difficult but sufficient to obtain good convergence in the elastic differential cross sections. We have also used pseudopotentials for the carbon and oxygen atoms. This strategy allows a reduction in the number of Cartesian Gaussian functions, since it is not necessary to consider those involved in the description of all $1s$ orbitals of these atoms.

Another way of calculating the scattering amplitude is by also expanding $\mathbf{k}_i$ in partial waves as follows:

$$\langle \mathbf{k}_f | f | \mathbf{k}_i \rangle = \sum_{l=0}^{\ell_{\max}} \sum_{m=-\ell}^{\ell} \sum_{l'=0}^{\ell_{\max}} \sum_{m'=-\ell'}^{\ell'} \langle \mathbf{k}_f | \ell m \rangle f(\ell m, \ell' m') \langle \ell' m' | \mathbf{k}_i \rangle. \quad (6)$$

The DCSs are obtained from $f(\ell m, \ell' m') = \langle \ell, m | f | \ell', m' \rangle$ which represents a scattering amplitude of an electron entering the interaction region in a partial wave $|\ell', m'\rangle$ and leaving it in a partial wave $|\ell, m\rangle$. The expansion up to $\ell_{\max} = 4$ is a standard procedure used in other scattering computer codes and, in order to allow future comparisons, we have included results calculated from Eq. (6) with several values of $\ell_{\max}$. For some cases (elastic and dipole-allowed singlet transitions), a Born-closure scheme was used following the same strategy as described in Ref. 11. This closure is obtained from the expression

$$f_{\text{LAB}}^{\text{closure}}(\mathbf{k}_f, \mathbf{k}_i) = f_{\text{LAB}}^{\text{FBA}}(\mathbf{k}_f, \mathbf{k}_i) + \sum_{\ell=0}^{\ell_{\max}} \sum_{m=-\ell}^{\ell} \left(f_{\text{LAB}}(\ell m, \mathbf{k}_i) - f_{\text{LAB}}^{\text{FBA}}(\ell m, \mathbf{k}_i) \right) Y_{\ell m}^*(\mathbf{k}_f), \quad (7)$$

where $f_{\text{LAB}}^{\text{FBA}}$ is the scattering amplitude for the permanent dipole moment potential for the elastic process or for the dipole transition potential for inelastic dipole-allowed processes. Both are obtained in the first Born approximation, in a closed form in the laboratory-frame. The amplitude $f_{\text{LAB}}(\ell m, \mathbf{k}_i)$ is just the $f(\ell m, \mathbf{k}_i)$ of Eqs. (5) or (6) transformed to this frame.

For the total cross section, we have used the optical theorem directly on the linear representation of the scattering amplitude, given by Eq. (1), i.e.,

$$\sigma_{\text{tot}} = \frac{1}{4\pi} \int d\Omega_{\mathbf{k}_i} \frac{4\pi}{k} \text{Im} f(\mathbf{k}_i, \mathbf{k}_i) = \sum_{n=1}^{N_{\text{open}}} \sigma(1 \rightarrow n), \quad (8)$$

where $\sigma(1 \rightarrow n)$ is the integral cross section (ICS) for the electronic transition $1 \rightarrow n$. In order to correct the contribution from the low-angular region (important for the elastic transition since phenol has a non-negligible permanent dipole moment, and for singlet excitations to dipole-allowed

states), we have also calculated the TCS, using the amplitudes obtained from Eq. (7), and the expression

$$\sigma_{\text{tot}}^{\text{closure}} = \sum_{n=1}^{N_{\text{open}}} \sigma^{\text{closure}}(1 \rightarrow n), \quad (9)$$

where $\sigma^{\text{closure}}(1 \rightarrow n) \neq \sigma(1 \rightarrow n)$ only for the above mentioned transitions.

We have also employed the binary-encounter-Bethe (BEB) model¹² to estimate the total ionization cross section by electron impact. This cross section, added to our computed elastic+inelastic cross section, will be compared to the total cross section data. This model provides the following expression for the ionization cross section per molecular orbital:

$$\sigma_{\text{BEB}}(T) = \frac{S}{t + u + 1} \left[\frac{\ln t}{2} \left(1 - \frac{1}{t^2} \right) + 1 - \frac{1}{t} - \frac{\ln t}{t + 1} \right]. \quad (10)$$

In Eq. (10), T is the incident electron energy, $t = T/B$ and $u = U/B$ are normalized energies, where B and U are the orbital binding and electron kinetic energy, respectively, and $S = 4\pi a_0^2 N_{\text{occ}} R^2 / B^2$, where N_{occ} is the orbital occupation number, $a_0 = 0.5292 \text{ \AA}$, and $R = 13.61 \text{ eV}$. The total ionization cross section is obtained by the summation of $\sigma_{\text{BEB}}(T)$ over the molecular orbitals satisfying $T > B$. This model provides ionization cross sections which agree with experiment to within 5%-15% (considering different molecules) for incident energies ranging from the first ionization threshold to several keV.¹³

III. COMPUTATIONAL PROCEDURES OF THE PRESENT APPLICATION OF THE SMCPP METHOD

In the present application of the SMCPP method, we have used the same Cartesian Gaussian basis set of Ref. 14 and the description of the target ground and excited states as in Ref. 15. This description can be summarized as (1) we have used 247 uncontracted Cartesian Gaussian functions for the description of target bound states and for the scattering calculations (CSFs basis of Eq. (4)), with s , p , and d types on the C and O atoms, and s and p types on the H atoms; (2) as usual,¹⁴ a Hartree-Fock calculation was carried out for the ground state, giving a permanent dipole moment of 1.423 D against the experimental value of 1.224 D of Ref. 16; (3) following the same strategy as used in our recent electron-ethylene multichannel study⁴ and as described in Ref. 15, a full single excitation configuration-interaction (FSCI) was obtained,¹⁷ so as to give excitation energies in reasonable agreement with the experimental spectra and time-dependent density functional theory calculations,¹⁵ particularly so for states lying below an arbitrary chosen energy of 7 eV; (4) a small set of singlet and triplet improved virtual orbitals were obtained and used as a minimum orbital basis for a single configuration interaction (MOBSCI)⁹ calculation, that was capable of describing all states below 7 eV in good agreement with the FSCI calculation. This was possible with 16 hole-particle pairs of single excitations, giving rise to 16 triplet

and 16 singlet states with 7 triplets and 5 singlets below our arbitrary chosen 7 eV energy.

The need for choosing a small value of energy (7 eV) is due to the high cost of the scattering calculation, when considering too big a value of N_{open} in Eq. (3) (for every electronic channel, a large numerical quadrature must be carried out in order to evaluate the projected Green's function¹⁸ of Eq. (2)). Our present CSF space for the scattering calculations has a smaller size than that used in our previous investigation¹⁴ and it was also constructed with modified virtual orbitals (in addition to the active particle orbitals of the MOBSCI) in order to better account for polarization effects. It is worthwhile noting that the configuration space is the same for all the present N_{open} ch-sep approximations (that is, for calculations with $N_{open} = 1, 2 \dots$ or 33, where more collision channels are opened as the electron impact energy increases). It is also important to mention that the present 1ch-sep cross sections obtained with this procedure are in close agreement with those previously reported,¹⁴ highlighting that the effect which will be discussed arises mainly from the different channel coupling schemes used in the present study. We have used the acronym 1ch-sep for the elastic process only, 3ch-sep for the ground plus first two triplet states (all states assigned to the elastic channel plus Band 1 of Ref. 15), the 8ch-sep for the ground state plus 5 triplets and two singlets (all states assigned to the elastic channel plus Bands 1 and 2 of Ref. 15), and 13ch-sep for the ground state plus 7 triplets and 5 singlets (all states assigned to the elastic channel plus Bands 1, 2, and 3 of Ref. 15). Other acronyms used in this paper are 21ch-sep, 23ch-sep, 28ch-sep, and 33ch-sep, meaning that all states of the MOBSCI calculation up to 10 eV, 15 eV, 20 eV, and 24 eV, respectively, are energetically accessible, and as a consequence, they are included in the sum appearing in the projector of Eq. (3). A summary of the theoretical spectrum is given in Table I. Note that the 33ch-sep approximation has all the MOBSCI channels (real and pseudostates) as electronically open channels. The

ionization, B , and kinetic, U , orbital energies necessary to obtain the BEB cross section were obtained in a Hartree-Fock calculation with the 6-311++G(2d,1p) basis set using the program GAMESS.¹⁷ The geometry employed in this calculation was optimized at the MP2 level with the same basis set. The calculated ionization potential of phenol is 8.61 eV which agrees well with the experimental value of 8.51 eV¹⁹ and the value observed in a recent electron impact ionization investigation.²⁰

IV. ATOMIC OPTICAL MODEL (OM)

The independent atom model with screening corrected additivity rule (IAM-SCAR) method has had significant use in calculating electron scattering cross sections for a wide variety of molecular targets (see, e.g., Refs. 21–26) and over a broad energy (E_0) range (typically ~ 1 –5000 eV). As a consequence, we only précis the key points of that approach here.

The first subjects of our computations are the individual atoms that form phenol, namely, carbon (C), oxygen (O), and hydrogen (H). The atomic OM is based on a potential scattering approach, where the local complex potential $V(r)$ is given by

$$V(r) = V_s(r) + V_{ex}(r) + V_p(r) + iV_a(r). \quad (11)$$

In the above equation $V_s(r)$ is the standard Hartree potential of the target, $V_{ex}(r)$ represents the exchange interaction of Riley and Truhlar,²⁷ $V_p(r)$ is the polarisation potential of Zhang *et al.*,²⁸ and $V_a(r)$ is the imaginary absorption potential of Staszewska *et al.*²⁹ Due to the last term in Eq. (1), the OM potential approach yields a complex phase shift. This allows for the calculation of the DCSs and ICSs for elastic and inelastic (taken to mean all excited electronic states and ionisation combined together) scattering, as well as the TCS as the sum of those ICSs for each atom.

TABLE I. Summary of the calculated and experimental excitation energies below 7 eV for phenol from Ref. 15. The MOBSCI calculation aimed to reproduce all 12 states of the FSCI calculation lying below 7 eV. These are 7 triplet and 5 singlet states listed below. The remaining MOBSCI pseudostates are 11 singlet states opening up at 8.31, 8.83, 8.89, 9.37, 9.45, 14.43, 17.55, 19.15, 20.34, 21.35, and 23.82 eV and 9 triplet states at 8.25, 8.37, 8.79, 12.65, 16.01, 18.46, 19.75, 21.10, and 23.26 eV.

	Expt. band	Symmetry	Dipole transition (D)	Energy (eV)		
				FSCI	MOBSCI	TD-DFT
TRIPLET	Band 1 (3.4–4.3 eV)	A'	...	3.29	3.57	3.71
	Band 1 (3.4–4.3 eV)	A'	...	4.49	4.73	4.10
	Band 2 (4.3–5.4 eV)	A'	...	4.78	4.90	4.53
	Band 2 (4.3–5.4 eV)	A''	...	5.94	6.03	5.06
	Band 2 (4.3–5.4 eV)	A'	...	5.73	6.16	5.30
	Band 3 (5.4–6.3 eV)	A''	...	6.53	6.78	5.53
	Band 3 (5.4–6.3 eV)	A''	...	6.86	6.99	5.90
SINGLET	Band 2 (4.3–5.4 eV)	A'	1.04	5.82	6.09	4.99
	Band 2 (4.3–5.4 eV)	A''	0.08	6.06	6.21	5.13
	Band 3 (5.4–6.3 eV)	A'	0.35	6.12	6.80	5.76
	Band 3 (5.4–6.3 eV)	A''	1.03	6.68	6.86	5.57
	Band 3 (5.4–6.3 eV)	A''	0.03	6.86	6.99	5.98

V. SCREENING CORRECTED ADDITIVITY RULE (SCAR) AND ROTATIONAL EXCITATIONS

We subsequently calculate the electron- $\text{C}_6\text{H}_5\text{OH}$ cross sections by applying the additivity rule (AR) to the OM results of each atom. In that approach, the molecular scattering amplitude stems from the sum of all the relevant atomic amplitudes, including the phase coefficients, which gives the elastic DCSs for phenol. Elastic ICSs can then be determined by integrating those DCSs. The sum of the elastic and absorption ICSs (for all inelastic processes except rotations and vibrations) then gives the molecular TCS. However, the AR does not account for the target molecular structure, so that it is only applicable when the incident electrons are so fast that they effectively “see” the target molecule as a sum of the individual atoms (typically above ~ 100 eV). To try and overcome this limitation, Blanco and García^{30,31} introduced the SCAR method which employs some screening coefficients to account for the geometry of the molecule (atomic positions and bond lengths). The IAM-SCAR method as described does not account for rotational and vibrational excitations. However, for polar molecules such as $\text{C}_6\text{H}_5\text{OH}$, additional dipole-induced excitation cross sections can be computed using the approach of Jain.³² In that method, rotational excitation DCSs and ICSs for a free electric dipole are calculated within the framework of the first Born approximation. Those results can now be incorporated into our IAM-SCAR results in an incoherent way, just by adding up the cross sections as independent channels. We call the cross sections that result from this latter process as “IAM-SCAR + rotations.”

There is a body of evidence that suggests that the IAM-SCAR approach for non-polar molecules and IAM-SCAR + rotations method for polar species^{21–26} provides a good description of the measured cross sections down to $E_0 = 20$ eV. However, there are also some examples^{33–35} where the comparison between these two approaches and the measured cross sections is only reasonable at energies above ~ 50 eV. Therefore, one of the rationales of the present investigation was to test the efficacy of the IAM-SCAR plus rotations results against corresponding results from our fully *ab initio* SMCPP at the N_{open} -ch-sep level computations, at representative energies of 50 eV and below.

VI. EXPERIMENTAL APPARATUS FOR THE TOTAL CROSS SECTIONS: DETAILS AND METHOD

The present total cross sections were measured with an apparatus, located in Madrid, whose functionality is based on a magnetic confinement of an electron beam when it enters into the collision chamber until its ultimate detection. The role of the main magnetic field (0.2 T) is simply to translate the electrons, which exhibit the precise angle and energy that resulted from a potential collision, along the central axis to the end of the collision chamber. This apparatus has already been recently described in detail,^{21,22} so that only a brief summary of its more important points is given here.

The apparatus consists of three regions: (i) electron gun region, (ii) collision chamber region, and (iii) the analyzer–detector region. Each region is connected by small

orifices and each has independent magnetic fields. Electrons are generated by thermionic emission from a tungsten filament and accelerated to a kinetic energy E_0 before passing into the collision chamber. The magnetic field of the electron gun region (B_g), oriented opposing the main field B , ensures a low angular divergence of this electron beam by locally compensating for B and preventing electrons leaving the filament in oblique directions to pass through the collimators. The collision chamber has a geometrical length of 140 mm and so is sufficiently large compared to the limiting apertures (of 1 mm and 2.3 mm diameter) to ensure a well-defined region of constant pressure. Phenol is introduced into the system using a heated variable leak valve, from a steel sample-container maintained at around 50 °C by means of various silicone heater mats. The chamber wall in principle can partly absorb the heat dissipated by the magnetic coils, depending on the pumping speed of the cooling liquid through the chamber’s water jacket.^{21,22} Hence, the balance between solenoid current and water speed can be utilized in order to maintain a convenient chamber temperature and avoid condensation of the phenol on the inner chamber walls. In this application, the collision cell temperature was held at the same temperature as the sample-container to ensure the phenol pressure was maintained at a constant value. The pressure in the scattering chamber (or cell) is determined by an absolute capacitance manometer (MKS Baratron 627B, Germany), while the temperature is measured (and monitored) using a K-type thermocouple in thermal contact with the inner chamber wall. Due to the magnetic confinement, the effective localization of the electrons after scattering (but before entering the analyzer/detector region) is within 1 mm around the central axis. After traversing the collision chamber, electrons are selected by a retarding potential analyzer (RPA) so that only those with parallel (axial) components of kinetic energy $\geq eV_R$ (where $e = 1.6 \times 10^{-19}$ C and V_R = retarding potential) continue towards the detector. The detector assembly is formed by two micro-channel plates and an anode, and is typically polarized with a voltage drop of around +2 kV. It is operated in a single-pulse counting mode and connected, via some additional electronics, to a computer running a custom LabView (National Instruments) program which registers and analyses the signal.

Phenol, with a stated purity better than 99%, is a white solid at room temperature but readily sublimates under vacuum. In this case, our sample was purchased from Sigma-Aldrich and it was further purified through the performance of freeze-pump-thaw cycles. Before each experiment, the energy resolution defined as $\Delta E = e(V_{R,90} - V_{R,10})/2$, with $V_{R,90}$ and $V_{R,10}$ being the retarding potentials leading to 90% and 10% of the transmitted electrons, respectively, was determined from the transmission curve $I(V_R)$ in vacuum. Here, I is the transmitted beam intensity (electron count rate). This energy resolution was usually found to be almost identical to the full-width-at-half-maximum (FWHM) of the derivative of $I(V_R)$ and is given, at each energy, in Table II. Subsequently, the retarding potential was fixed to define the cut-off energy^{21,22} and then the transmitted intensity as a function of the gas pressure was recorded. In the present study, the pressure range was ~ 0.2 –4 mTorr, depending on the energy, but always ensuring the attenuation was less than 50% of the incident

TABLE II. Measured total cross sections (σ_{exp}), the correction to the TCSs due to the forward angle scattering effect (σ_{forw}) and the “true” TCSs (σ), as a function of energy, for electron collisions with phenol, in 10^{-16} cm^2 units. For each energy, the reproducibility index, energy resolution (ΔE) and angular discrimination ($\Delta\theta$) are also given.

E_0 (eV)	σ_{exp} (10^{-16} cm^2)	Reproducibility (%)	ΔE (eV)	$\Delta\theta$ ($^\circ$)	σ_{forw} (10^{-16} cm^2)	σ (10^{-16} cm^2)
10	34.8	12	3.5	36.3	18.0	52.8
15	36.8	7	1.8	20.3	12.3	49.1
20	30.6	8	2.2	19.4	11.5	42.1
30	30.1	5	2.0	15.0	7.72	37.8
40	31.2	8	2.3	13.9	6.95	38.2
50	29.9	12	1.8	10.9	5.43	35.2
70	25.1	8	2.8	11.5	4.88	30.0
100	23.7	2	2.2	8.5	4.15	27.9
150	18.5	6	2.0	6.6	3.12	21.6
200	18.1	6	2.2	6.0	2.58	20.7
300	14.5	6	1.9	4.6	2.04	16.5
500	9.7	6	2.7	4.2	1.51	11.2
700	8.7	4	2.3	3.3	1.24	9.9
1000	5.9	3	2.4	2.8	0.97	6.9

intensity in order to avoid multiple scattering effects. For each incident electron energy in the range of 10–1000 eV (see Table II), a series of 7–10 attenuation curves, each normally comprising 7–12 data pairs (I and pressure (p)), were then acquired. Those data points were now fitted with an exponential curve $I(p)$, from which the experimental TCS, σ_{exp} , is obtained according to the Beer-Lambert law

$$I = I_0 e^{-nl\sigma_{\text{exp}}} = I_0 e^{-pl\sigma_{\text{exp}}/kT}. \quad (12)$$

Note that in Eq. (12), I_0 is the intensity of the in vacuum non-attenuated beam, n is the number density of the phenol gas, $l = 0.1413 \text{ m}$ is the effective collision chamber length,²¹ k is Boltzmann’s constant, and $T = \sqrt{T_c T_m}$ is the gas temperature (K) calculated according to the thermal transpiration effect³⁶ between the manometer at T_m and the collision chamber at T_c .

The experimental uncertainties on our TCSs are no worse than $\sim 10\%$, including contributions due to uncertainties in the collision chamber length, the phenol gas pressure (the major contributor), the temperature, and the true incident beam energy. Furthermore a “reproducibility index,” reflecting the standard deviation between sets of $I(p)$ versus p curves from the same series, in the range of 2%–12% was also found (see Table II for the exact value at each E_0). The “reproducibility index” comprised contributions from the filament emission stability, temperature stability, (electronic) signal fluctuations, and an uncertainty associated with the fitting to the attenuation curves. Combining the aforementioned factors, one obtains a general precision in the TCS determination in the range of 8%–16%, depending on the energy, for our 10–1000 eV investigation.

In addition to this largely statistical uncertainty, the angular acceptance $\Delta\theta$ of the apparatus is a limiting aspect and represents a potentially important source of systematic error. Here, elastically or rotationally and vibrationally inelastically scattered electrons at small angles ($0^\circ \leq \theta \leq \Delta\theta^\circ$) or near 180° ($180^\circ - \Delta\theta^\circ \leq \theta \leq 180^\circ$), for the case of backscattered electrons, are not distinguished from the unscattered electrons

and so will not be accounted for in our determination of the measured TCS (σ_{exp}). Under these circumstances, σ_{exp} will always be lower (and sometimes significantly lower²¹) than the “true” value of the TCS (σ). The angular discrimination in the forward direction ($\Delta\theta$) can be determined from the energy resolution ΔE via^{21,22}

$$\Delta\theta = \sin^{-1}(\sqrt{\Delta E/E_0}). \quad (13)$$

The values of $\Delta\theta$ for the present measurements were found to lie between 3° and 36° , and are listed in Table II. This means that any valid comparison to our SMCPP + BEB and IAM-SCAR + rotations theory TCSs needs to take into account that the present measured values σ_{exp} actually represent only partial values

$$\sigma_{\text{exp}}(E) \approx \sigma(E) - \sigma_{\text{forw}}(E). \quad (14)$$

Our procedure for calculating $\sigma_{\text{forw}}(E)$, using the IAM-SCAR elastic DCS and the Born-based rotational DCS, the vibrational scattering cross sections being ignored as they are expected to be relatively much smaller in magnitude in our 10–1000 eV energy range,³⁷ can be found in Fuss *et al.*²¹ and Sanz *et al.*²² and so we do not repeat that detail here. Rather, we simply note that our values for σ_{forw} and σ can be found in Table II. We reiterate only a comparison of $\sigma(E)$ with corresponding theory is meaningful and we discuss this in more detail shortly in our relevant results and discussion section.

VII. RESULTS AND DISCUSSION: ELECTRONIC MULTICHANNEL EFFECTS ON THE ELASTIC ELECTRON-PHENOL DIFFERENTIAL CROSS SECTIONS

Figures 1–3 show the ℓ_{max} convergence for our calculated differential cross sections for a selection of phenol’s excited electronic states, obtained with the scattering amplitude of Eq. (6), with respect to the converged result obtained from

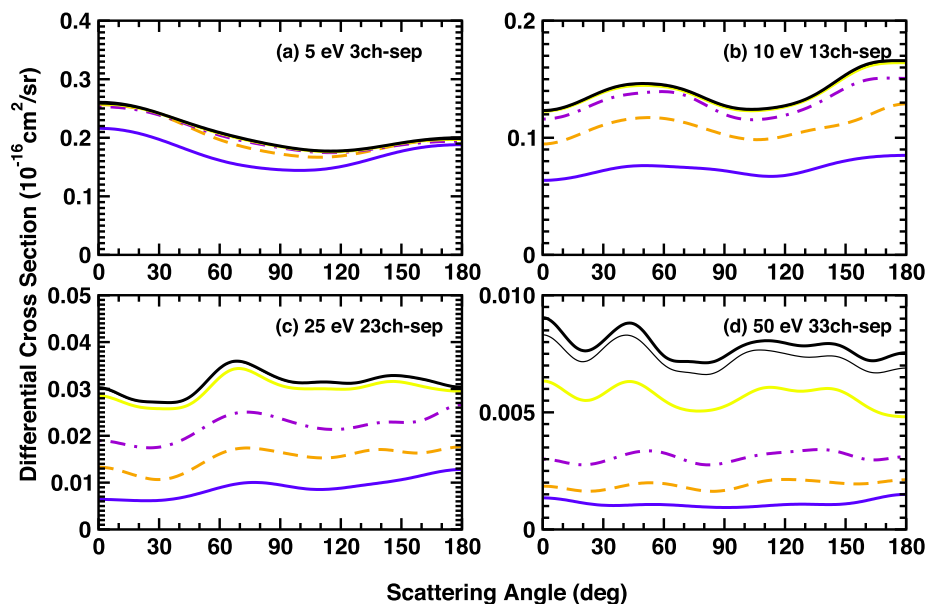


FIG. 1. Differential cross sections obtained from the scattering amplitude given by Eq. (6) for the excitation of the first triplet state of phenol at (a) 5.0 eV at a 3ch-sep approximation, (b) 10 eV at a 13ch-sep approximation, (c) 25 eV at a 23ch-sep approximation, and (d) 50 eV at a 33ch-sep approximation. Full (blue) line: $\ell_{\max} = 4$; long dashed (orange) line: $\ell_{\max} = 5$; dashed-dotted (violet) line: $\ell_{\max} = 6$; full (yellow) line: $\ell_{\max} = 8$; thin full (black) line: $\ell_{\max} = 10$; thick full (black) line: converged result from Eq. (5) with $\ell_{\max} = 13$.

Eq. (5), with $\ell_{\max} = 13$. We have performed this study for one excited state from each experimental band (1st triplet state of phenol that belongs to Band 1, 1st singlet of phenol that belongs to Band 2, and 4th singlet state of phenol that belongs to Band 3) of the experimentally resolved bands in the energy loss spectra of Ref. 15. The electronic excitation cross sections of those bands are the subject of study of an accompanying paper⁵ in this same journal issue. Figure 1 shows the DCSs for 5, 10, 25, and 50 eV for the 1st triplet of phenol calculated at the 3ch-sep, 13ch-sep, 23ch-sep, and 33ch-sep levels of approximation, respectively. In Figure 2, we present our DCSs for 8, 15, 25, and 50 eV for the 1st singlet of phenol calculated at the 8ch-sep, 13ch-sep, 23ch-sep, and 33ch-sep levels of approximation, respectively. Next, in Figure 3, we show DCSs for 8, 15, 25, and 50 eV for the 4th singlet of phenol (chosen because it has the largest dipole transition of Band 3¹⁵) calculated at the 8ch-sep, 23ch-sep, 28ch-sep, and 33ch-sep levels of approximation, respectively.

The results of these 3 figures clearly show that $\ell_{\max} = 4$ is not enough for convergence, even for an impact energy as low as 5 eV or 8 eV for the first triplet or first singlet, respectively. For a reasonable convergence, $\ell_{\max} = 8$ is necessary for energies up to 25 eV, and for higher energies, $\ell_{\max} = 10$ or 13 is required. For 50 eV, the DCS obtained with the scattering amplitude from Eq. (6) with $\ell_{\max} = 4$ can be less than 10% of the converged DCS. This is clearly unsatisfactory. For justifying the use of $\ell_{\max} = 4$ in Eq. (6), we would need an indication that the higher partial wave contributions from the SMCPP method are less significant than simply zeroing them out (for the triplet transitions) or less significant than replacing them with a 1st Born approximation for the higher partial waves (for singlet transitions with the Born closure). Unfortunately, as will be seen in the elastic study ahead, there is no indication that we could cut $\ell_{\max}$ to a value as low as 4.

Figures 4–6 show the multichannel effects on the 1st triplet state, and on the selected¹⁵ Band 2 of phenol for impact

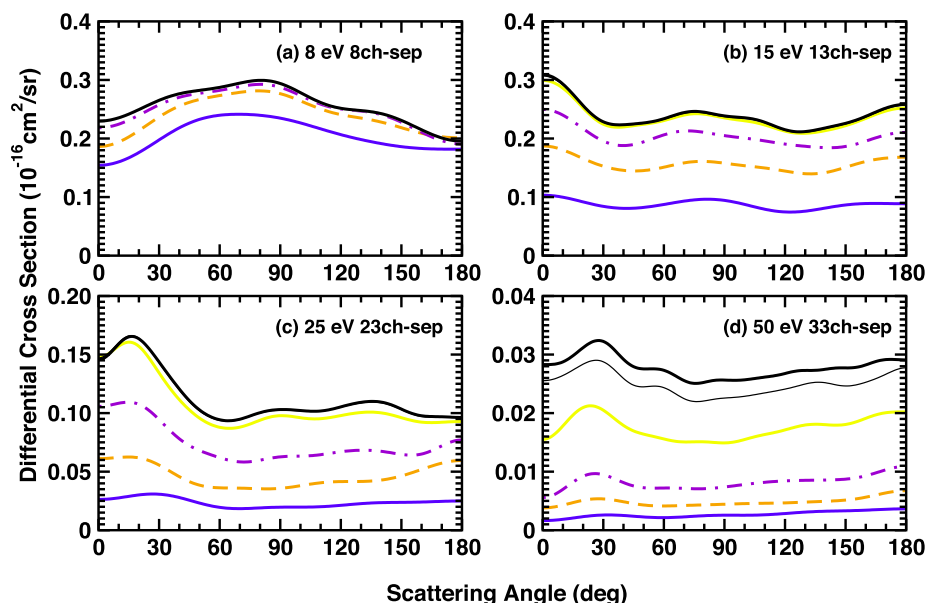


FIG. 2. Differential cross sections obtained from the scattering amplitude given by Eq. (6) (no closure) for the excitation of the first singlet state of phenol. Labels for $\ell_{\max}$ are the same as those used in Fig. 1, but at (a), 8ch-sep level of approximation for 8 eV, (b) 13ch-sep for 15 eV, (c) 23ch-sep for 25 eV, and (d) 33ch-sep for 50 eV electron impact energy.

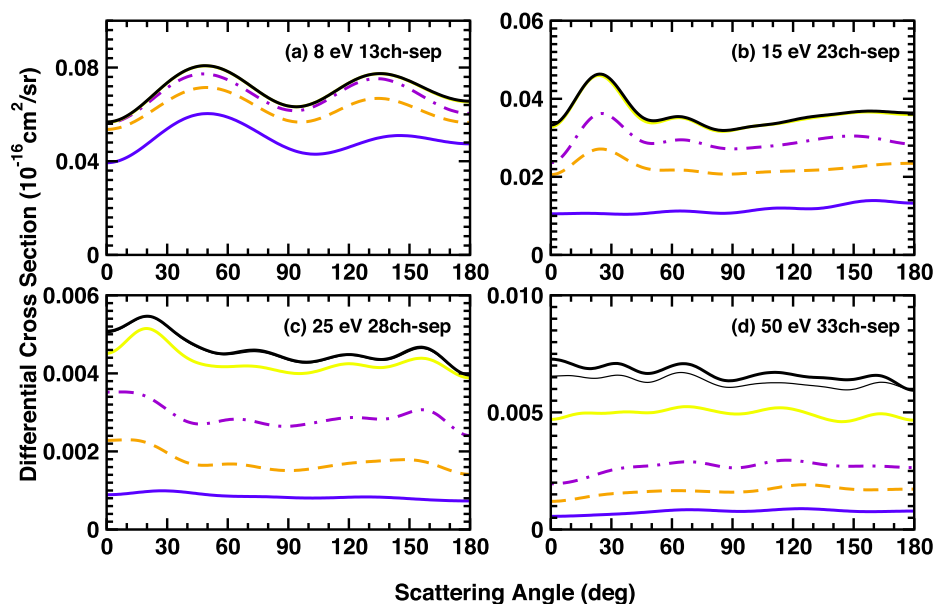


FIG. 3. Differential cross sections obtained from the scattering amplitude given by Eq. (6) (no closure) for the excitation of the fourth singlet state of phenol. Labels for ℓ_{max} are the same as those used in Fig. 1, but at (a), 13ch-sep level of approximation for 8 eV, (b) 23ch-sep for 15 eV, (c) 28ch-sep for 25 eV, and (d) 33ch-sep for 50 eV electron impact energy.

energies varying from 8 to 50 eV. In particular, Figure 4 shows these effects on the electronic excitation DCSs of the $^3A'$ state of phenol (see Table I) for 8, 10, 12, and 15 eV, at the 3ch-sep, 8ch-sep, 13ch-sep, 21ch-sep, and 23ch-sep levels. The level of coupling is defined by the electronic states that are energetically open according to the impact energy and thresholds given in Table I. Figure 5 shows the multichannel influence at the higher impact energies, 20, 30, 40, and 50 eV for this transition. The channel couplings shown in this figure are at the 3ch-sep, 8ch-sep, 13ch-sep, 23ch-sep, 28ch-sep, and the 33ch-sep levels of approximation. In accord with our previous study with ethylene,⁴ we see for phenol that by increasing the number of channels, the probability flux to a particular state ($^3A'$ in this case) decreases as we open more electronic states, especially at the higher energies. About the same trend is also seen in Fig. 6 for the electronic excitation of the experimental Band 2, formed by 3 triplets and 2 singlets (see Table I), for electron impact energy ranging again from 8

to 50 eV. In general, the results follow the trend of decreasing the flux into the relevant band excitation as we open more channels. In this case, however, there are a few exceptions, e.g., at 15 eV electron impact energy, the 13ch-sep DCSs appear stronger in magnitude than the 8ch-sep results and there is also an inversion between the 28ch-sep and 33ch-sep DCSs for 30 and 50 eV. Nonetheless, we reiterate that the overall trend we have highlighted typically holds for the states incorporated in Band 2. Further discussion on the electronic excitation of phenol is carried out in an accompanying paper⁵ also published in this journal issue.

Figures 7 and 8 show elastic DCSs and the influence of multichannel effects in this transition. As in the inelastic case, the DCSs decrease their magnitude as we increase the number of open channels. We compare our elastic phenol differential cross sections with the experimental data obtained by Cho *et al.*³⁸ for benzene, a molecule differing from phenol by the replacement of a H atom in benzene with a Hydroxyl

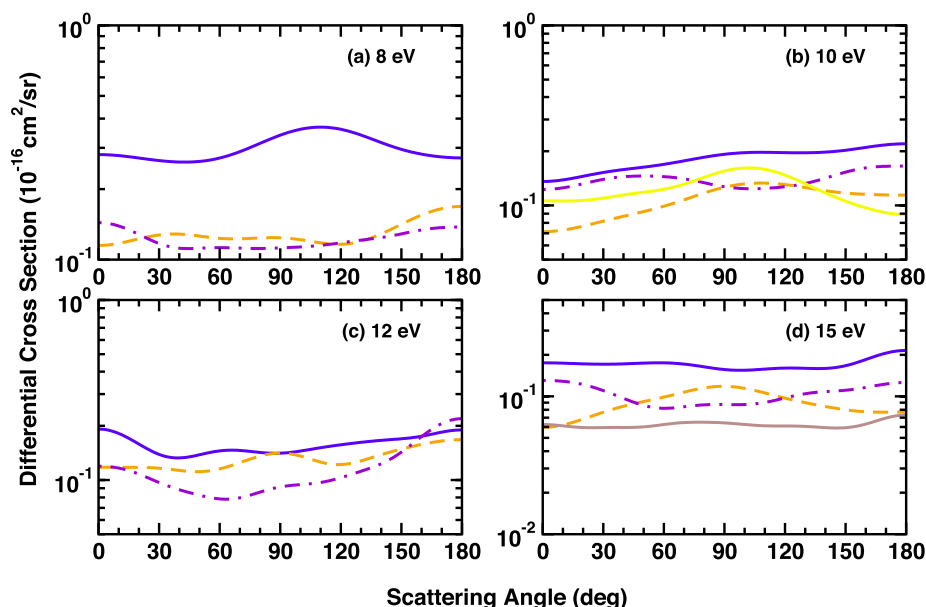


FIG. 4. Multichannel coupling influence on the differential cross sections for the excitation of the first triplet state of phenol at (a) 8.0 eV, (b) 10.0 eV, (c) 12.0 eV, and (d) 15.0 eV. Full (blue) line: 3ch-sep; dashed (orange) line: 8ch-sep; dashed-dotted (violet) line: 13ch-sep; full (yellow) line: 21ch-sep; full (brown) line: 23ch-sep.

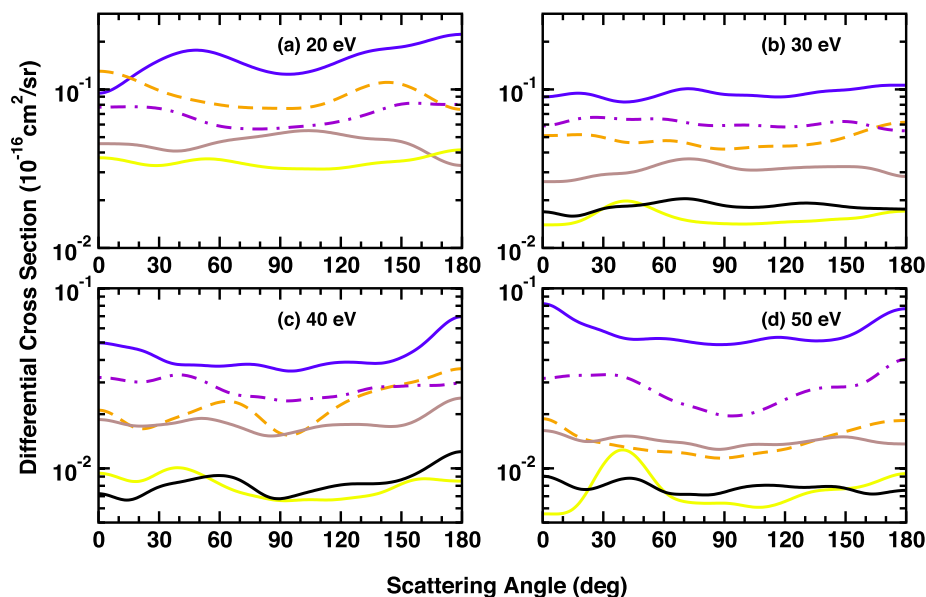


FIG. 5. Multichannel coupling influence on the differential cross sections for the excitation of the first triplet state of phenol at (a) 20.0 eV, (b) 30.0 eV, (c) 40.0 eV, and (d) 50.0 eV. Full (blue) line: 3ch-sep; dashed (orange) line: 8ch-sep; dashed-dotted (violet) line: 13ch-sep; full (brown) line: 23ch-sep; full (yellow) line: 28ch-sep; full (black) line: 33ch-sep.

(OH) group. Benzene has no dipole moment, but as a result of this replacement, phenol molecules have a permanent dipole moment. The dipole moment in phenol is not too large (66% of the water molecule dipole) and affects mainly the scattering cross sections in the more forward direction. Figure 7 shows elastic DCSs for impact energies of 5, 8, 10, and 12 eV. It is possible to note from this figure that at 5 eV, the multichannel effect seems to be not too important. On the other hand, for higher energies, the presence of multichannel effects changes substantially the DCS when compared to the one obtained at a 1ch-sep level of approximation. At 8 eV, the 3ch-sep, 8ch-sep, and 13ch-sep elastic cross sections are in better agreement with the experiments than the usual sep approximation which usually has only the elastic channel as open (here, named 1ch-sep). To our surprise, at 10 eV, the 13ch-sep approximation gave worse results than the 3ch-sep and 8ch-sep calculations when compared with the experimental data for benzene. However, as we further increase the number of channels to a

21ch-sep level of approximation, the DCS for 10 eV increases in size and gets in much better agreement with the experiments. The nature of this instability for some of the N_{open} -ch-sep approximations will be discussed in Sec. VIII. Figure 8 shows the present calculated elastic DCSs for impact energies at 15, 20, 30, and 40 eV. The agreement with the experiments is excellent for all energies and, in general, follows the trend that as we open more channels the scattered elastic flux decreases towards the experimental data of benzene. This level of agreement in the benzene and phenol elastic DCSs is somewhat intriguing, as they are different species (although both have ring structures), and we believe might reflect the fact, at least in part, that they both possess very similar dipole polarizabilities.¹⁶ Such a good agreement invites us to speculate that the inelastic cross section should also match well with our new experiments, which is not the case, as discussed in detail in the accompanying inelastic electronic-state paper.⁵

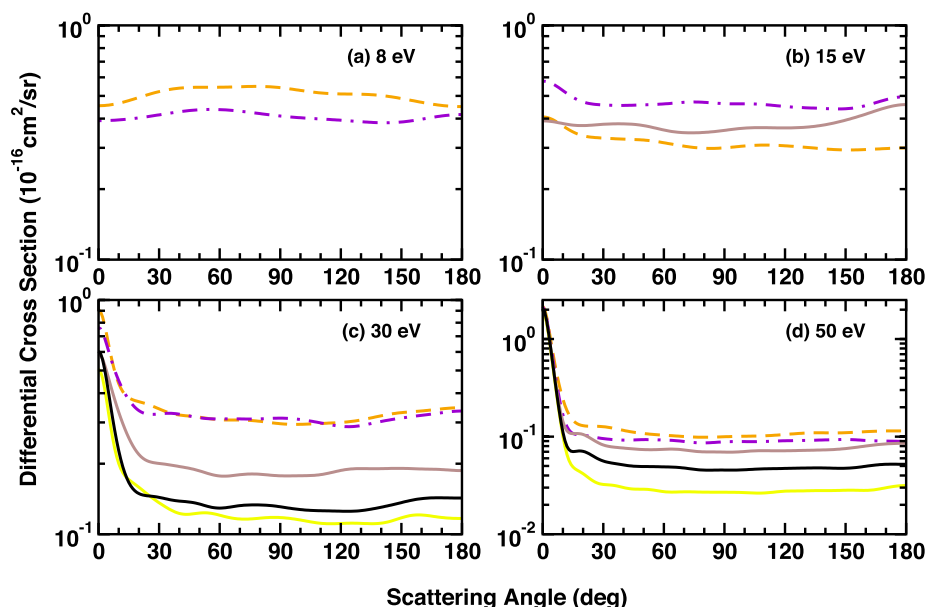


FIG. 6. Multichannel coupling influence on the differential cross sections obtained from the scattering amplitude given by Eq. (7) (with closure, $\ell_{max} = 10$) for the excitation of our selected¹⁵ Band 2 of phenol at (a) 8 eV, (b) 15 eV, (c) 30 eV, and (d) 50 eV. Dashed (orange) line: 8ch-sep; dashed-dotted (violet) line: 13ch-sep; full (brown) line: 23ch-sep; full (yellow) line: 28ch-sep; full (black) line: 33ch-sep.

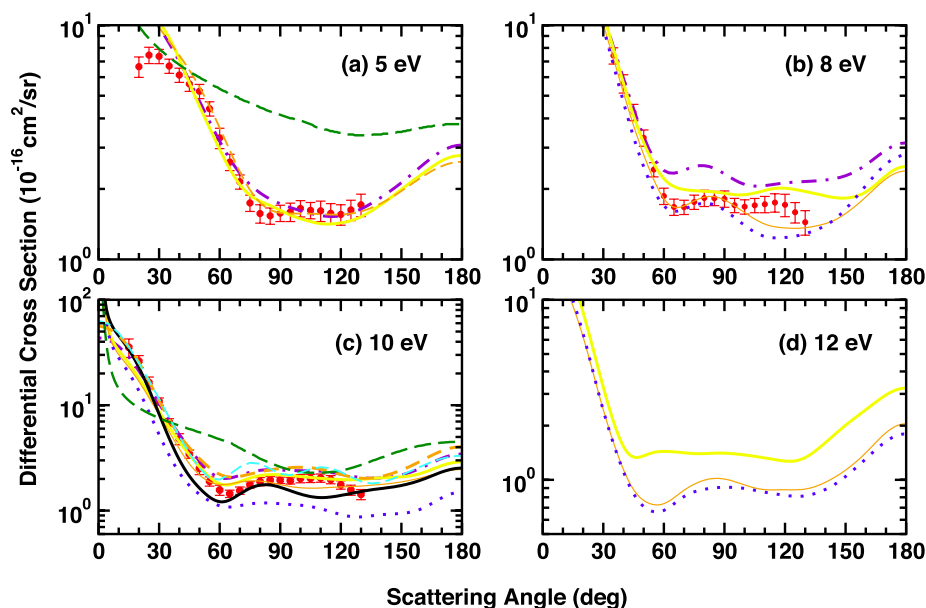


FIG. 7. Multichannel coupling influence on the elastic differential cross sections of phenol for (a) 5.0 eV, (b) 8.0 eV, (c) 10.0 eV, and (d) 12.0 eV electron impact energies. Full (red) circles: experiments of Ref. 38 for benzene; dashed (orange) line: 1ch-sep of Ref. 14; dashed-dotted (violet) line: 1ch-sep; full (yellow) line: 3ch-sep; thin full (orange) line: 8ch-sep; dotted (blue) line: 13ch-sep; full (black) line: 21ch-sep; thin dashed (green): present IAM-SCAR calculation; dashed (cyan) line: SVIM calculation for benzene of Ref. 39.

In terms of a comparison between our computed IAM-SCAR elastic DCS and the measured elastic DCS, we find that the level of overall accord is not as impressive as that just described for the converged SMCPP results. While our IAM-SCAR angular distributions do reproduce the gross features of the experimental angular distributions, particularly at the higher impact energies, some disagreement in terms of the absolute magnitude of the cross section are found in Figures 7 and 8. This is not a unique case, the IAM-SCAR approach had, at the lower energies (<50 eV), some difficulties in describing the elastic DCS for the radical CF_3 ,^{33,34} and the bio-molecule 1,4-dioxane³⁵ which we believe can be ascribed to the screening corrections not fulfilling their role as well as they might. Nonetheless, for phenol, the IAM-SCAR results are in quite fair agreement with the benzene DCS above about 40 eV energy. A better degree of agreement was found between the elastic DCSs calculated by the SMCPP for phenol

and those calculated by the Schwinger variational iterative method (SVIM) for benzene³⁹ at 10, 20, 30, and 40 eV. The SVIM is a single center expansion method that allows the calculation of cross sections via an optical potential model, which accounts for the contribution of an average number of electronically inelastic processes. Again, we see that the agreement is improved as the impact energy increases.

Figure 9 depicts the angular momentum $\ell_{\max}$ convergence, used in Eq. (5), to obtain the elastic differential cross sections. This figure shows that $\ell_{\max} = 4$ can be used only for low energies such as 5 eV. As the energy increases, the use of such a low value of angular momentum gives an artificially “wavy” DCS which is smaller than the converged results in magnitude. The good agreement with the measured benzene DCSs, seen in the previous figures, indicates that at least for elastic transitions we should not rely on such small $\ell_{\max}$ values.

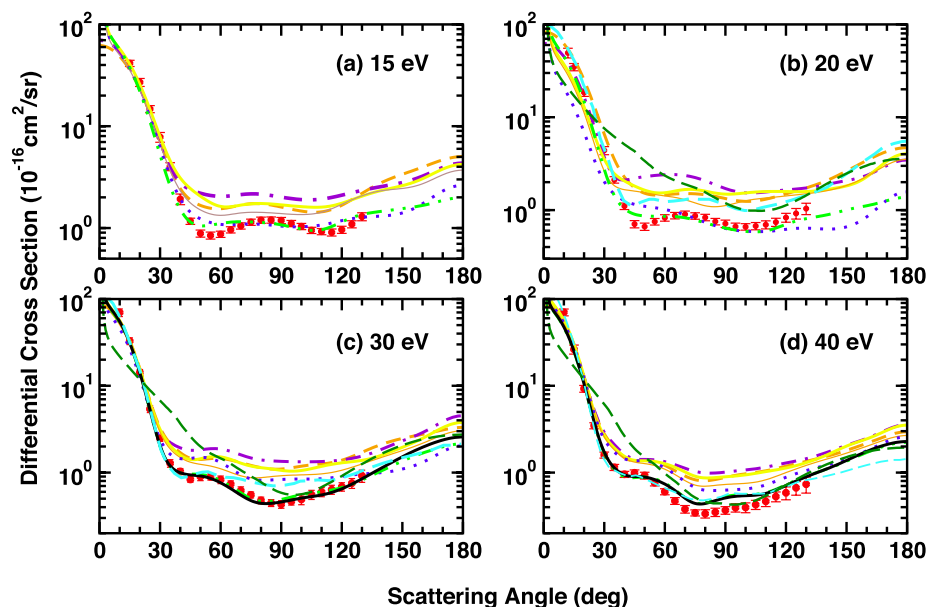


FIG. 8. Multichannel coupling influence on the elastic differential cross sections of phenol for (a) 15.0 eV, (b) 20.0 eV, (c) 30.0 eV, and (d) 40.0 eV electron impact energies. Full (red) circles: experiments of Ref. 38 for benzene; dashed (orange) line: 1ch-sep of Ref. 14 (unpublished DCS results); dashed-dotted (violet) line: 1ch-sep; full (yellow) line: 3ch-sep; thin full (orange) line: 8ch-sep; dotted (blue) line: 13ch-sep; dashed-dotted (green) line: 23ch-sep; thick full (black) line: 33ch-sep; thin dashed (green): present IAM-SCAR calculation; thick dashed (cyan) line: SVIM calculation for benzene of Ref. 39.

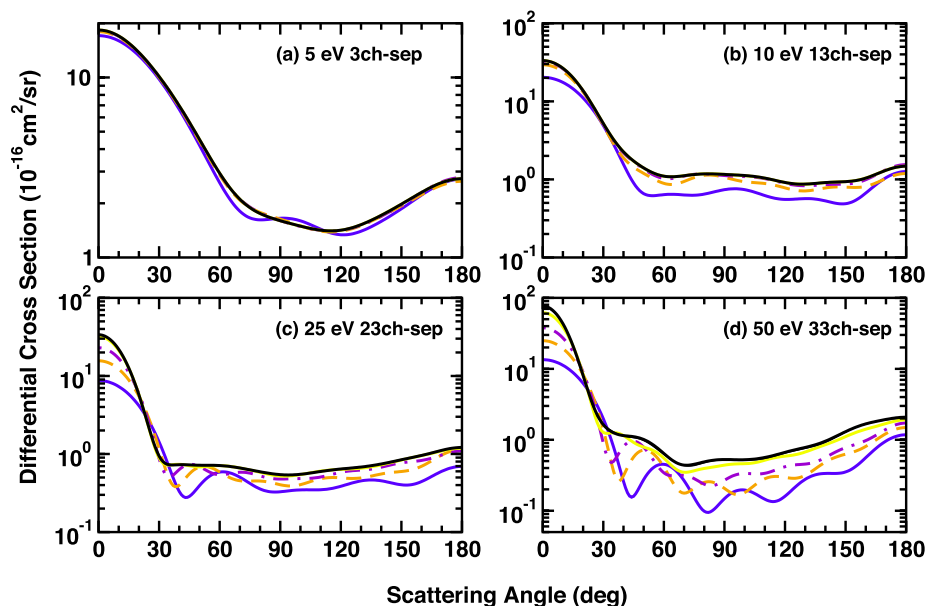


FIG. 9. Differential cross sections for elastic electron scattering by phenol with the same definitions as in Fig. 1.

VIII. RESULTS AND DISCUSSION: ELECTRON-PHENOL TOTAL, ELASTIC INTEGRAL, AND ELASTIC MOMENTUM TRANSFER CROSS SECTIONS

At low energies, only the elastic channel is electronically open and the distortion of the electronic cloud is taken into account by virtual excitations (closed channel space) from the ground state. However, as we increase the impact energy, other electronic states can be excited and if we do not open these channels (i.e., include them in the projector P of Eq. (3) for out-coming flux via the Green function), spurious structures (bumps and troughs) show up in the calculated cross section. These features are more clear in the integral, momentum transfer, and total cross sections, as we will show in the next figures. Due to the presence of $(1 - \cos \theta)$ in the angular integration to obtain the MTCS, a permanent dipole moment, even if it is large (resulting in a strong forward peaking DCS), does not contribute as much as in the calculation of the ICS. So, in principle, we would expect to have better agreement between phenol and benzene for the MTCS than for the ICS. However, as noted earlier, the dipole moment of phenol is relatively small and, as we will soon see, the elastic ICS is also comparing well with the corresponding results for benzene. Figure 10 brings a comparison between our calculated momentum transfer cross sections at several levels of approximations (1ch-sep, 3ch-sep, 13ch-sep, 23ch-sep, 28ch-sep, and 33ch-sep) and the experimental measurements of Cho *et al.*³⁸ for benzene. The connected up-triangles are the 1ch-sep calculation. The black curve is a combination of our best channel coupling results for each interval of energy, i.e., 1ch-sep for 1-4 eV, 3ch-sep for around 5 eV, 13ch-sep for 8-14 eV, 23ch-sep for around 15 eV, 28ch-sep for around 20 eV, and 33ch-sep for energies above 24.9 eV. Here, it is worthwhile noting that the best channel coupling is determined by two parameters: the energy of the incident electron and the thresholds of the excited states obtained according to the MOBSCI strategy. So, for example, for the incident electron energy of 8 eV, there are 13-states (the 7 triplets and the 5

singlets displayed in Table I plus the ground state) that became energetically allowed. All other 20 excited states, referred to in the text as pseudostates (for details, please, see the caption of Table I), have their thresholds above 8 eV and were treated as closed channels in this case. Our calculated MTCSs show clearly the trend seen before, i.e., the scattered flux decreases in the elastic channel as a result of opening more electronic channels. The 23ch-sep for 15 and 20 eV (not as much), and the 33ch-sep approximations for 30 and 40 eV are in very reasonable agreement with the experiments. At lower energy, the 13ch-sep approximation gives good results for 8 and 10 eV, as we should expect from the comparison of the DCS results. Spikes and troughs are present in this figure, and in order

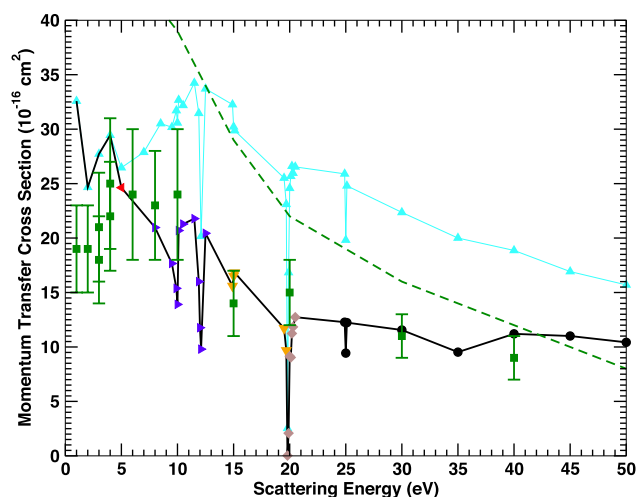


FIG. 10. MTCSs for phenol in the energy region 0-50 eV. Up (cyan) triangles: 1ch-sep; left (red) triangles: 3ch-sep; right (blue) triangles: 13ch-sep; down (orange) triangles: 23ch-sep; (brown) diamonds: 28ch-sep; filled (black) circles: 33ch-sep; full (black) curve: connecting our best SMCPP channel coupling results for each interval of energy (see text); thin dashed (green): present IAM-SCAR calculation; filled (dark green) squares: experimental results of Ref. 38 for benzene. SMCPP results presented in this figure were obtained in a calculation without use of the Born-closure procedure (see Sec. IX for details).

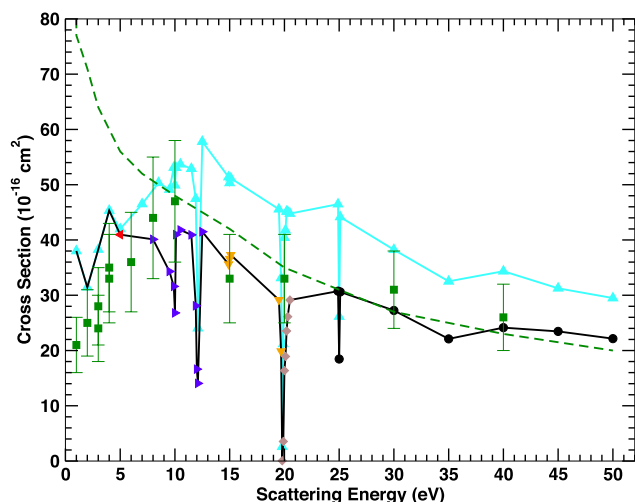


FIG. 11. ICS for phenol in the energy region 0-50 eV. Up (cyan) triangles: 1ch-sep; full (black) curve: our best SMCPP channel coupling results for each interval of energy (see text for every energy interval we used the same symbols as in Fig. 10); thin dashed (green): present IAM-SCAR calculation; filled (dark green) squares: experimental results of Ref. 38 for benzene. SMCPP results presented in this figure were obtained in a calculation without use of the Born-closure procedure (see Sec. IX for details).

to properly identify the spurious features from the physical features, we needed a larger number of energy points. At 25 eV, we could see that the calculations are quite unstable going from (13ch-sep or 28ch-sep) to (23ch-sep or 33ch-sep). By narrowing the grid of energies, we found a strong and narrow trough around 25 eV, due to some spurious closed channel in our CSF space that should in fact be open. Similar features are also seen at 10, 12, and 20 eV impact energies. The latter is due to the threshold of a triplet pseudostate state that opens at 19.75 eV (see Table I). Other narrow spikes and troughs could be found, and in order to properly identify the spurious features from the physical features, we would need a much larger number of points, which would make the calculation very expensive. A similar study (same symbols) is shown in Fig. 11 for our elastic ICS in comparison with the experimental data of Cho *et al.*³⁸ for benzene. Above 3 eV, the integral cross sections with and without closure are very close to each other (consistent with phenol's small dipole moment) and only the latter are included for comparison with the experiments for benzene. The agreement here between our theoretical and the experimental elastic ICS is as good as the agreement found for the MTCS.

Figure 12 presents our TCS for electron-phenol scattering from 5 to 50 eV along with the experimental benzene data of Sueoka⁴⁰ and Mozejko *et al.*⁴¹ Our SMCPP cross section when summed with the ionization cross section obtained with the BEB model is in very good agreement with present TCS experimental data. The agreement between the SMCPP TCSs obtained with 33ch-sep, full (black) line, and 1ch-sep, full (cyan) line connecting up-triangles, reveals that the latter cannot give good elastic cross sections. As we have discussed previously in the ethylene study,⁴ the present application also shows that as we open more channels, the flux is distributed among them, decreasing substantially the elastic cross sections (clearly noted in our DCSs figures). The 33ch-sep SMCPP

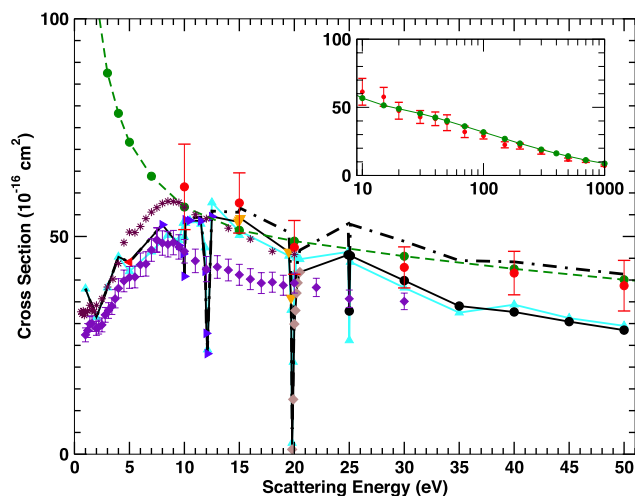


FIG. 12. TCS for phenol in the energy region 0-50 eV. The SMCPP calculations were obtained from Eq. (8). Up (cyan) triangles: 1ch-sep; full (black) line: connecting our best SMCPP channel coupling results for each interval of energy (see text for every energy interval we used the same symbols as in Figs. 10 and 11); thick dotted-dashed line: our best SMCPP TCS plus total ionization cross sections using the BEB approximation; thin dashed line with filled (green) circles: TCS of the IAM-SCAR method; filled (violet) diamonds, (brown) stars, and filled (red) circles: results of the benzene TCS experiments of Ref. 40, of Ref. 41, and present phenol TCS data. The inset shows our theoretical TCS obtained with the IAM-SCAR method and our present experimental phenol TCS data for higher energies. SMCPP results presented in this figure were obtained in a calculation without use of the Born-closure procedure (see Sec. IX for details).

TCS was obtained with the optical theorem, Eq. (8), without closure. If we add the ICSs, channel by channel, as in Eq. (9), and include the BEB result then the resulting TCS is less than 3% larger than the one obtained by Eq. (8) for energies above 5 eV. This is a nice level of self-consistency from the two approaches to derive a SMCPP-based TCS. At very low energies (below 3 eV), our SMCPP-based TCS does not present a strong increase, shown by the present IAM-SCAR TCSs, which is characteristic of a polar molecule since we have not included the contribution due to the long range dipole potential through the Born closure. The present IAM-SCAR calculated TCS is also in good accord with our measured phenol data, to within the uncertainties on the data, and over the common energy range between 10 and 50 eV. Included as an inset of Fig. 12 is a plot showing the higher-energy TCS behavior. As expected, our measured TCS phenol data are in good accord with our IAM-SCAR results at the higher energies.

IX. A LITTLE COMMENT ON THE CHOICE OF ℓ_{max}

Since the issue of convergence of the cross sections in terms of the parameter ℓ_{max} permeates all discussion of the results presented above, we believe that it is appropriate to introduce a complementary comment in order to put the subject on a more general perspective. The Born-closure affects basically the forward scattering. Therefore, for scattering angles above 30°–40°, the SMCPP+closure DCSs would be approximately equal to the SMCPP itself (and comparable to experiments). So, Figures 2, 3, and 9 would continue showing convergent results only for $\ell_{max} = 10$ even if we applied the

Born-closure scheme. We feel confident in applying the Born-closure strategy when, by incrementing the value of ℓ_{max} , only the forward scattering of the SMCPP DCS is affected. For elastic DCSs, this is not the case even for 10 eV, as can be seen in Figure 9. In order to be sure that the Born-closure scheme with $\ell_{max} = 4$ could be the right choice, we need an argument that only the higher partial waves of the SMCPP are wrong and therefore deserve to be deleted and replaced by the first Born approximation. Another possibility would be that as we open more and more channels, the lower partial wave contributions would converge and the higher partial wave contribution would decrease. Unfortunately, the whole spectra of partial wave contributions decrease all together with the flux distribution. Therefore, considering the excellent agreement with the experimental data for Benzene in the elastic channel (with multichannel effects and $\ell_{max} = 10$), we decided to publish only the DCSs obtained with the converged (backward scattering) SMCPP results, which produces ICS, MTCS, and TCS very little affected by the Born-closure procedure.

X. CONCLUSIONS

We have reported experimental and theoretical cross sections for scattering of low-energy electrons by phenol. The SMCPP method was employed to calculate elastic differential, integral, momentum transfer, and total cross sections according to the so-called N_{open} -ch-sep level of approximation. More specifically, up to thirty three electronic states of phenol, obtained within a single excitation configuration-interaction description, were included in the open-channel space. Through the use of this channel coupling scheme, we have performed a detailed analysis of the influence of multichannel coupling effects on the DCS results. As a general trend, we found that the probability flux into the elastic channel decreases as more channels were included in the calculation, leading to a substantial reduction on the magnitude of the corresponding cross sections. This effect becomes even more important for higher impact energies, as expected. The IAM-SCAR method was also used to compute DCSs for energies above 10 eV, and the agreement with the SMCPP results was improved as the energy increased. The SMCPP cross sections obtained were in very good agreement with the experimental data measured by Cho *et al.* for benzene,³⁸ both in terms of magnitude and angular distributions.

Integral and momentum transfer cross sections for elastic electron scattering by phenol, obtained with the SMCPP and the IAM-SCAR methods, were compared to the measurements performed for benzene by Cho *et al.*³⁸ The overall good agreement with that experimental data clearly indicates that the inclusion of multichannel coupling effects (from the 3ch-sep up to the 33ch-sep level of approximation) was essential for an adequate description of the scattering process under consideration, especially if compared to the cross sections (DCSs, MTCSs and ICSs) obtained at the 1ch-sep level of approximation, where only the elastic channel was included in the scattering calculation.

The present total cross sections in the energy range from 10 to 50 eV included measurements as well as calculated

results obtained by using the IAM-SCAR and the SMCPP (sum of elastic+inelastic+ionization contributions) methods. All these results are in good agreement with each other within the uncertainties associated to the experimental data. In addition, our measured data and IAM-SCAR results at higher energies were also found to be in excellent accord.

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