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Deviations from reversible dynamics in a qubit–oscillator system coupled to a very small environment



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HIGHLIGHTS

- Model considering two qubits coupled to an oscillator.
- One of the qubits (qubit2) acts as a small environment.
- Reduced evolution of the other qubit (qubit1)—quasi-periodic behaviour.
- Linear entropy of qubit1 is very sensitive to the degree of mixedness of qubit2.

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ABSTRACT

In this contribution it is considered a simple and solvable model consisting of a qubit in interaction with an oscillator exposed to a very small “environment” (a second qubit). An isolated qubit–oscillator system having the oscillator initially in one of its energy eigenstates exhibits Rabi oscillations, an evidence of coherent quantum behaviour. It is shown here in which way the coupling to a small “environment” disrupts such regular behaviour, leading to a quasi-periodic dynamics for the qubit linear entropy. In particular, it is found that the linear entropy is very sensitive to the amount of mixedness of the “environment”. For completeness, fluctuations in the oscillator energy are also taken into account.

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1. Introduction

Exactly solvable models often give new insights into the properties of simple quantum systems. A well known of such models is the Jaynes–Cummings model [1] of quantum optics: a two-level system (qubit) coupled to the quantized field (oscillator) under the rotating wave approximation (RWA). A peculiar feature of that model is the periodic exchange of excitations between the qubit and the oscillator, namely Rabi oscillations of the qubit population, if the oscillator is initially prepared in a number state $|n\rangle$, an energy eigenstate of the free oscillator. A similar behaviour is found for other properties of the qubit, such as the von Neumann entropy, for instance, a quantifier of the purity of the qubit state. Perfect Rabi oscillations of the entropy mean that a system initially in a pure state is able to return to a pure state at a later time, and they are therefore a signature of coherent quantum behaviour. On the other hand, if the oscillator is initially prepared in a distribution of number states, e.g., a coherent state, entropy will be no longer a periodic function of time, and the qubit will not be able to return to a pure state [2,3]. Thus, if the qubit is coupled to a quantum oscillator in a state which is itself a superposition of number states, the qubit evolution, in particular, its entropy, will become significantly disordered. This is because the incommensurate Rabi frequencies $\Omega \propto \sqrt{n+1}$ present in the Jaynes–Cummings solution introduce dephasing, leading to

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an irreversible-like behaviour. Another possible source of disorder, even for the oscillator initially in a single number state $|n\rangle$ is related to chaos, and arises in a strongly coupled qubit–oscillator system without making the RWA [4].

In addition to that, the investigation of the evolution of simple quantum systems coupled to an environment has been a subject of continuing interest for many years [5]. In very precise experiments, it has been found that the Rabi oscillations in a Jaynes–Cummings system are indeed irreversibly damped [6]. This important phenomenon is attributed to external influences (noise), which are normally modelled by coupling the system of interest to a very large number of other quantum systems (reservoir). Those phenomenological models [7], based on the assumption that a large system having certain pre-defined properties is coupled to a subsystem of interest, are in general able to describe the main features of dissipative quantum subsystems: loss of energy and loss of quantum coherence. However, rather than being coupled to a large reservoir, a quantum system may be subjected to an uncontrollable coupling to external systems having few degrees of freedom, and one may ask in which ways such a small “environment” may disturb the quantum features of the system of interest. In this paper I am going to consider a model consisting of a two-level system (qubit1) interacting with an oscillator weakly coupled to a single sub-system constituted by a second two-level system (qubit2) which will in fact play the role of a small “environment”. This corresponds to another well known model of quantum optics: the two-atom Tavis–Cummings model [8,9]. Here, though, I will explore this simple quantum system from a different perspective: the qubit1–oscillator will be our system of interest, and qubit2 will constitute a perturbing environment. One of the advantages of employing the Tavis–Cummings model, is that it is exactly solvable under the RWA, which will allow a general, non-perturbative treatment of the dynamics of qubit1. The main purpose of this contribution is to find in which way quantum reversibility represented by the Rabi oscillations of entropy in the number-state Jaynes–Cummings model could be affected by coupling the oscillator to a small system and how that disturbance depends on the features of such an “environment”. In particular, I am going to analyse the dynamics of the linear entropy of qubit1 for different degrees of mixedness of the “environment” (qubit2). I would like to point out that it is not my intention to present a model to reproduce a phenomenon such as irreversible loss of quantum coherence, but rather, I am interested in examining the deviations from perfect reversibility due to a (unwanted) coupling to a “small environment”.

The paper is organized as follows. In Section 2, I will present a preliminary discussion based on the Jaynes–Cummings model for simple preparations of the oscillator (pure and mixed). In Section 3, I will show the solution of the actual system–“environment” model and the calculation of the linear entropy of qubit1. In Section 4, I will discuss the time-evolution of the linear entropy for different initial conditions as well as the influence of energy fluctuations in the oscillator. In Section 5 I will summarize the conclusions.

2. Jaynes–Cummings model with an initial mixed state

In the “pure state” Jaynes–Cummings model, if the oscillator is initially prepared in a number state, the entropy has a regular behaviour; perfect Rabi oscillations, as just a single frequency is involved. However, if the oscillator is initially in a mixed state, different excitation numbers will come into play, yielding different Rabi frequencies. A simple situation considering a small number of excitations for the oscillator could be an initial mixed state involving just the vacuum $|0\rangle$ and the first excited state $|1\rangle$

$$\rho_{\text{osc}}(0) = f|0\rangle\langle 0| + (1-f)|1\rangle\langle 1|. \quad (1)$$

Here qubit1 (with states $|e\rangle$ and $|g\rangle$) is assumed to be initially in the excited state $|e\rangle$. The qubit1–oscillator joint density operator (Jaynes–Cummings) may then be written as

$$\rho(t) = \hat{A}(t)\rho_o(0)\hat{A}^\dagger(t)|g\rangle\langle g| + \hat{A}(t)\rho_o(0)\hat{B}^\dagger(t)|g\rangle\langle e| + \hat{B}(t)\rho_o(0)\hat{A}^\dagger(t)|e\rangle\langle g| + \hat{B}(t)\rho_o(0)\hat{B}^\dagger(t)|e\rangle\langle e|, \quad (2)$$

where $\hat{A}(t) = \cos(\lambda t \sqrt{\hat{a}\hat{a}^\dagger})$ and $\hat{B}(t) = -i\hat{a}^\dagger \sin(\lambda t \sqrt{\hat{a}\hat{a}^\dagger})/\sqrt{\hat{a}\hat{a}^\dagger}$. Here, a (a^\dagger) are the oscillator annihilation (creation) operators and λ is the qubit–oscillator coupling constant. After tracing over the oscillator’s degrees of freedom, one obtains the reduced density operator for qubit1, $\rho_{q1}(t) = \text{Tr}_o[\rho(t)]$. Then, a straightforward calculation gives us the linear entropy¹ for qubit1, $\zeta(t) = 1 - \text{Tr}[\rho_{q1}^2(t)]$, or

$$\zeta(t) = 1 - [f \cos^2(\lambda t) + (1-f) \cos^2(\sqrt{2}\lambda t)]^2 - [f \sin^2(\lambda t) + (1-f) \sin^2(\sqrt{2}\lambda t)]^2. \quad (3)$$

That is an example of a “quasi periodic function”, a sum of periodic functions with incommensurate frequencies $\nu_1 = \lambda$ and $\nu_2 = \sqrt{2}\lambda$. In Fig. 1 it is plotted ζ as a function of time (normalized to $\lambda = 1.0$) for $f = 0.5$. We note a damping of the oscillations, followed by an increase of their amplitude so that ζ gets very close to zero again after a few cycles. Strictly speaking, though, qubit1 does not return to a pure state, although ζ may get very close to zero in this case. This simple example will be useful to understand the behaviour of the qubit1–oscillator dynamics coupled to an external system, as follows.

¹ The linear entropy is zero for a pure state ($\text{Tr}\rho^2 = \text{Tr}\rho = 1$), and different from zero for a statistical mixture. In our specific case, the linear entropy varies from 0 to 0.5.

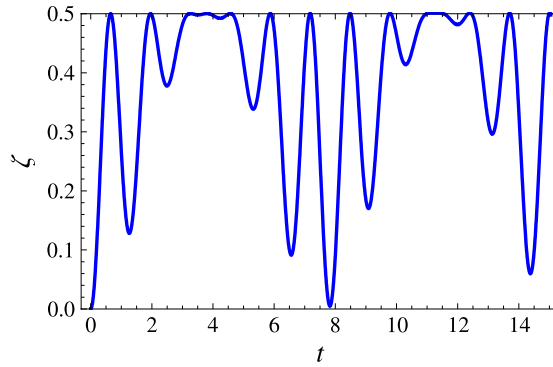


Fig. 1. Plot of the linear entropy ζ (as a function of t) of a qubit initially in state $|e\rangle$ and the oscillator initially in the mixed state $\rho_{\text{osc}}(0) = f|0\rangle\langle 0| + (1 - f)|1\rangle\langle 1|$ with $\lambda = 1.0$ and $f = 0.5$.

3. The model

The model considered here consists of a two-level system, qubit1, coupled to a quantized field mode (oscillator), which is itself coupled to a second two-level system: qubit2, the “environment”. The Hamiltonian describing such a system (within the RWA) is the familiar two-atom Tavis–Cummings Hamiltonian [8,9], or

$$H = \hbar\omega_0 a^\dagger a + \frac{\hbar}{2} \sum_{i=1}^2 \omega_{i,q} \sigma_i^z + \hbar \sum_{i=1}^2 \lambda_i (a \sigma_i^+ + a^\dagger \sigma_i^-), \quad (4)$$

where a (a^\dagger) are the oscillator annihilation (creation) operators, and σ_i^z , σ_i^+ , σ_i^- are the Pauli operators for the i th qubit. I am going to consider the fully resonant case, i.e., having both qubit1 and qubit2 frequencies equal to the oscillator frequency, $\omega_{i,q} = \omega_0$. However, the oscillator–qubits couplings, λ_1 and λ_2 will be assumed to be different, so that it will be straightforward to simulate a weakly coupled subsystem. That model has a well known exact analytical solution [9], but here I am going to use initial conditions different than the usual ones: while qubit1 and the oscillator will be assumed to be initially prepared in pure states, qubit2, the “environment”, will be in a more general, mixed state. The initial state of the whole system will be given by the product state $\rho(0) = \rho_1(0) \otimes \rho_{\text{osc}}(0) \otimes \rho_2(0)$, where

$$\rho_1(0) = |e_1\rangle\langle e_1|, \quad \rho_{\text{osc}}(0) = |\phi\rangle\langle\phi|, \quad \rho_2(0) = p|e_2\rangle\langle e_2| + (1 - p)|g_2\rangle\langle g_2|, \quad (5)$$

or qubit1 in the upper state, the oscillator in a generic pure state $|\phi\rangle = \sum B_n |n\rangle$, (with B_n real) and qubit2 in a statistical mixture of upper and lower states. The time-evolved system density operator may be expressed as a sum of two terms

$$\rho(t) = p |\Psi_p(t)\rangle\langle\Psi_p(t)| + (1 - p) |\Psi_{1-p}(t)\rangle\langle\Psi_{1-p}(t)|, \quad (6)$$

where

$$|\Psi_p(t)\rangle = \sum_n B_n \left[C_1(t) |e_1, e_2, n\rangle + C_2(t) |e_1, g_2, n+1\rangle + C_3(t) |g_1, e_2, n+1\rangle + C_4(t) |g_1, g_2, n+2\rangle \right], \quad (7)$$

and

$$|\Psi_{1-p}(t)\rangle = \sum_n B_n \left[C'_1(t) |e_1, e_2, n-1\rangle + C'_2(t) |e_1, g_2, n\rangle + C'_3(t) |g_1, e_2, n\rangle + C'_4(t) |g_1, g_2, n+1\rangle \right]. \quad (8)$$

After solving the Schrödinger equation, we obtain the coefficients $C(C')$, which are given by

$$\begin{aligned} C_1(t) &= \left(\frac{1}{2D} \right) a_- \cos d_+ t + a_+ \cos d_- t \\ C_2(t) &= - \left(\frac{i}{2D} \right) (a_- b_+)^{(1/2)} \sin d_+ t - (a_+ b_-)^{(1/2)} \sin d_- t \\ C_3(t) &= - \left(\frac{i}{2D} \right) (a_- b_-)^{(1/2)} \sin d_+ t - (a_+ b_+)^{(1/2)} \sin d_- t \\ C_4(t) &= \left(\frac{1}{2D} \right) (a_+ a_-)^{(1/2)} [\cos d_+ t - \cos d_- t], \end{aligned} \quad (9)$$

where $a_{\pm} = D \pm (\lambda_1^2 + \lambda_2^2)$, $b_{\pm} = D \pm (\lambda_1^2 - \lambda_2^2)$, with

$$D = [(2n+3)^2(\lambda_1^2 + \lambda_2^2)^2 - 4(n+1)(n+2)(\lambda_1^2 - \lambda_2^2)^2]^{1/2}, \quad (10)$$

$$d_{\pm} = [(2n+3)(\lambda_1 + \lambda_2)^2 \pm D]^{1/2} / \sqrt{2}. \quad (11)$$

Also

$$\begin{aligned} C'_1(t) &= -\left(\frac{i}{2D'}\right) (a'_- b'_+)^{(1/2)} \sin d'_+ t - (a'_+ b'_-)^{(1/2)} \sin d'_- t \\ C'_2(t) &= \left(\frac{1}{2D'}\right) b'_- \cos d'_+ t - b'_+ \cos d'_- t \\ C'_3(t) &= \left(\frac{1}{2D'}\right) (b'_+ b'_-)^{(1/2)} [\cos d'_+ t - \cos d'_- t] \\ C'_4(t) &= -\left(\frac{i}{2D'}\right) (a'_+ b'_+)^{(1/2)} \sin d'_+ t + (a'_- b'_-)^{(1/2)} \sin d'_- t. \end{aligned} \quad (12)$$

The “primes” in the expressions above (D' , etc.) mean that the corresponding terms are evaluated by replacing n by $n-1$.

I am primarily interested in investigating the influence of the environment (qubit2) on the dynamics of qubit1. For instance, in order to find deviations from reversible behaviour in the qubit1–oscillator interaction, we may calculate the linear entropy relative to qubit1, obtained from its reduced density operator by tracing over qubit2 and the oscillator variables, or $\rho_{q1}(t) = \text{Tr}_{q2,o} [\rho(t)]$. After some algebra, we obtain

$$\zeta = 1 - \text{Tr} [\rho_{q1}^2(t)] = 1 - (\alpha^2 + \beta^2 + 2|\gamma|^2), \quad (13)$$

where

$$\alpha = \sum_n (B_n)^2 [p (|C_{3,n}|^2 + |C_{4,n}|^2) + (1-p) (|C'_{3,n}|^2 + |C'_{4,n}|^2)], \quad (14)$$

$$\beta = \sum_n (B_n)^2 [p (|C_{1,n}|^2 + |C_{2,n}|^2) + (1-p) (|C'_{1,n}|^2 + |C'_{2,n}|^2)], \quad (15)$$

and

$$\gamma = \sum_n B_n B_{n+1} \left\{ p [C_{4,n} C_{2,n+1}^* + C_{3,n} C_{1,n+1}^*] + (1-p) [C'_{4,n} C'_{2,n+1}^* + C'_{3,n} C'_{1,n+1}^*] \right\}. \quad (16)$$

4. Linear entropy—analysis and numerical results

4.1. Oscillator initially prepared in a number state $|N\rangle$

Now I am going to analyse the time evolution of the linear entropy of qubit1 for different conditions of the oscillator and environment. I would like to recall that we are actually dealing here with systems having a few degrees of freedom, in contrast to the usual theory of open quantum systems, in which the relevant subsystems are assumed to be coupled to reservoirs with a large number of degrees of freedom. Our “environment” is constituted by one of the smallest possible disturbing systems: a single two-level system. As we have already seen, in the absence of an environment on any kind, if the oscillator is initially prepared in a number state (oscillator energy eigenstate), the linear entropy of qubit1 will be simply a periodic function of time, i.e., it will perform Rabi oscillations. This makes evident the perfectly reversible character of the qubit1–oscillator dynamics in such an ideal case. Now, what would be the influence of the small environment (qubit2) on the time evolution of the linear entropy of qubit1? Next I am going to analyse the behaviour of the linear entropy for different degrees of mixedness of the environment. The oscillator will be initially prepared in a number state $|N\rangle$, or $B_n = \delta_{n,N}$.

Firstly, I would like to point out the peculiar structure of Eq. (11); d_+ and d_- are basically the relevant frequencies of the dynamics, as the linear entropy contains terms oscillating at them. A very important point is that d_+ and d_- are in general incommensurate. Of course if the oscillator is completely decoupled from qubit2 ($\lambda_2 = 0.0$) we will have $d_+ = d_-$ and the linear entropy of qubit1 will be a truly periodic function. Otherwise, if we “turn on” the environment, the linear entropy will become a quasi periodic function of time, instead. We now assume a relatively weak coupling; for instance, by choosing $\lambda_2 = 0.1$ and $\lambda_1 = 10\lambda_2$. We may first consider the environment initially in a pure state ($p = 0.0$). In this case frequencies d_+ and d_- are different, the oscillations will be modulated, and we have again a quasi periodic function. In Fig. 2(a) we have a plot of the linear entropy as a function of time in that situation. The linear entropy increases but almost returns to zero after a few cycles, similarly to the case we have seen in Section 2, i.e., the oscillator initially in mixed state (see Fig. 1) but without the presence of an “environment”. This means that the effect (on qubit1) of an “environment” of this type (qubit2) initially

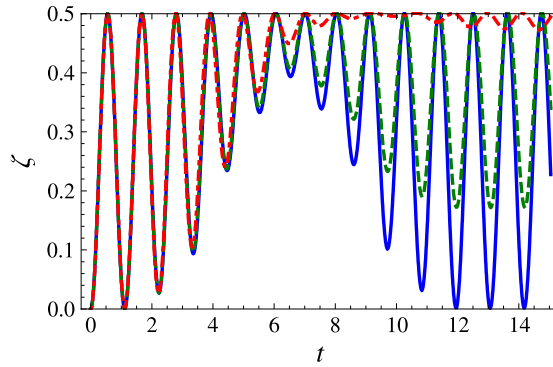


Fig. 2. Plot of the linear entropy ζ (as a function of t) of qubit1 initially in state $|e\rangle$, the oscillator initially in the number state $|1\rangle$ and qubit2 initially: (a) in a pure state, $p = 0.0$ (blue, continuous curve); (b) in a mixed state, $p = 0.1$ (green, dashed curve); (c) in a maximally mixed state, $p = 0.5$ (red, dot-dashed curve). In all cases $\lambda_1 = 1.0$ and $\lambda_2 = 0.1$. Note that all three curves virtually coincide for times $t < 5$.

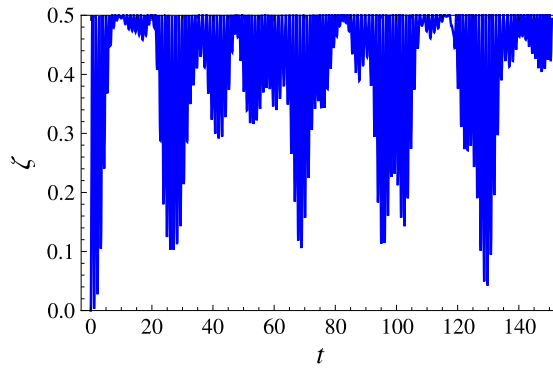


Fig. 3. Plot of the linear entropy ζ (as a function of t) of qubit1 initially in state $|e\rangle$, the oscillator initially in the number state $|1\rangle$ and qubit2 initially in a maximally mixed state, $p = 0.5$ for a longer time-scale, with $\lambda_1 = 1.0$ and $\lambda_2 = 0.1$.

in a pure state, may have some resemblance to the case in which there is no environment but the oscillator is initially in a mixed state.

However, the situation is quite different if the environment is initially in a mixed state; in this case terms containing two more different (and incommensurate) frequencies, d'_+ and d'_- , are introduced; see expressions (13)–(16). We expect, then, that this will cause a stronger dephasing. In fact, even a small amount of “mixedness” in qubit2 will bring a significant disorder to qubit1, as shown in Fig. 2(b), where qubit2 is in a mixed state with $p = 0.1$. Of course the effect of the “environment” is even more destructive if qubit2 is initially in a maximally mixed state, i.e., for $p = 0.5$. As shown in Fig. 2(c), in this case the linear entropy remains very close to its maximum value ($\zeta_{\max} = 0.5$) during a time interval corresponding to several cycles (after $t > 7$). In the discussion above a relatively short time-scale has been considered, but as a quasi periodic function is involved, we expect some sort of recurrence at longer times. Nevertheless, even for a larger time-scale, we note a very irregular evolution of ζ , indicating that the system spends most of the time in a highly mixed state, as shown in Fig. 3. As we would expect, a noisier environment (maximally mixed state) has a much stronger impact on qubit1 linear entropy ζ compared to a pure state environment case. Notwithstanding, this model does not lead to full irreversibility in the sense that the entropy grows to a maximum value for long times. Moreover, dissipation does not take place, as excitations go back and forth between system and environment.

4.2. Oscillator prepared in a binomial state: effect of energy fluctuations

Now suppose that in addition to the coupling of the oscillator to an external system (qubit2), there is some noise due to imperfections in the preparation of the state's oscillator, e.g., energy fluctuations. In order to take into account such effects, I will consider the oscillator initially prepared in a binomial state² [10], with coefficients B_n given by

$$B_n^M = \left[\frac{M!}{(M-n)!n!} q^n (1-q)^{M-n} \right]^{1/2}. \quad (17)$$

² The binomial state is characterized by two independent parameters: the maximum permissible number of excitations, M , and the probability of having a single excitation, q .

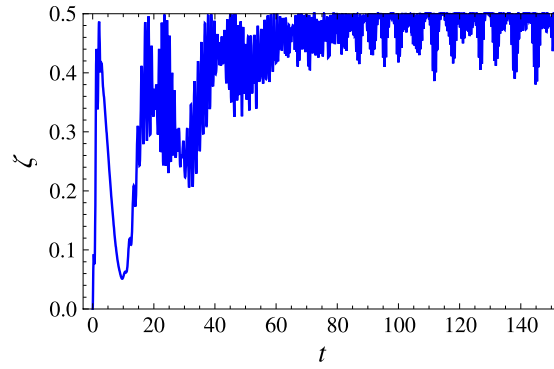


Fig. 4. Plot of the linear entropy ζ (as a function of t) of qubit1 initially in state $|e\rangle$ and the oscillator initially in a binomial state with $M = 100$ and $q = 0.1$. In this case qubit2 is decoupled, $\lambda_2 = 0.0$ and $\lambda_1 = 1.0$.

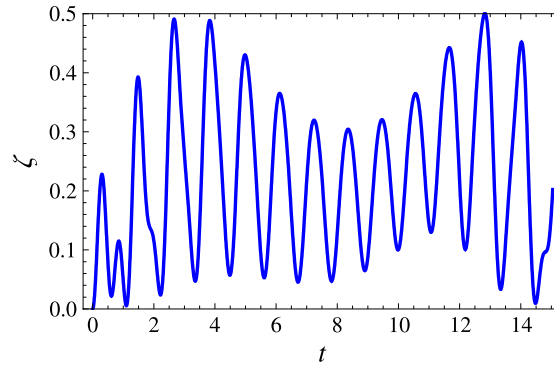


Fig. 5. Plot of the linear entropy ζ (as a function of t) of qubit1 initially in state $|e\rangle$ and the oscillator initially in a binomial state with $M = 7$ and $q = 0.85$. In this case qubit2 is decoupled, $\lambda_2 = 0.0$ and $\lambda_1 = 1.0$.

Therefore, in Eqs. (7) and (8) the summations are up to the finite number M , i.e., $|\phi\rangle = \sum_{n=0}^M B_n |n\rangle$. In Eqs. (14) and (15) the summations run up to $n = M$ and in (16) to $n = M - 1$. The binomial states have been already considered in the context of the two-atom Tavis–Cummings model [11], but the authors were concerned with the dynamics of atomic populations and photon number distributions, instead. Binomial states are in fact very convenient for my purposes mainly because of their interpolating properties: if the parameter $q = 1$, the binomial state becomes the number state $|M\rangle$; if $q \rightarrow 0$ and $M \rightarrow \infty$, but having $Mq = \alpha^2$ finite, the binomial state becomes the coherent state $|\alpha\rangle$. The energy fluctuations may be quantified via the Fano factor, defined as the dispersion of excitation number divided by the mean excitation number, or $F = \langle \Delta \hat{n}^2 \rangle / \langle \hat{n} \rangle$. For a binomial state, $F = 1 - q$. Note that $F = 0$ for a number state. First, I would like to make a preliminary discussion about the effect of fluctuations without having the oscillator coupled to qubit2.

4.2.1. Absence of an environment

It is well known that if the oscillator is initially prepared in a distribution of number states, rather than in number states, the von Neumann entropy (or the linear entropy) of the qubit will have a very different behaviour. For instance, if the oscillator is initially prepared in a (pure) coherent state, the qubit1 entropy will sharply increase to a maximum value and then decrease to a minimum that may be very close to zero [2,3]; a larger initial oscillator mean energy will bring the oscillator closer to a pure state at a specific time. As time goes on, the entropy grows again and starts showing an irregular behaviour. A very similar behaviour is shown in Fig. 4, for an initial binomial state having $M = 100$ and $q = 0.1$, i.e., a good approximation for a coherent state. In this case the entropy oscillates but it will grow progressively, showing that the system will remain in a mixed state even at very long times [3] (irreversible-like behaviour).

As an example of deviations from coherent oscillations characteristic of the number state case, we have in Fig. 5 the short time behaviour of the linear entropy for qubit1 in the case of a binomial state with $M = 7$ and $q = 0.85$. In this case the linear entropy also oscillates and at a later time it gets very close to zero. Therefore, a small energy spread in the initial oscillator state may also induce a quasi-periodic behaviour to the qubit1 linear entropy.

4.2.2. Presence of an environment

Now I would like to comment about the combined effects of fluctuations in excitations and the environment, as both favour irregular evolution of the linear entropy. Suppose that the oscillator has a relatively small energy spread, e.g., having

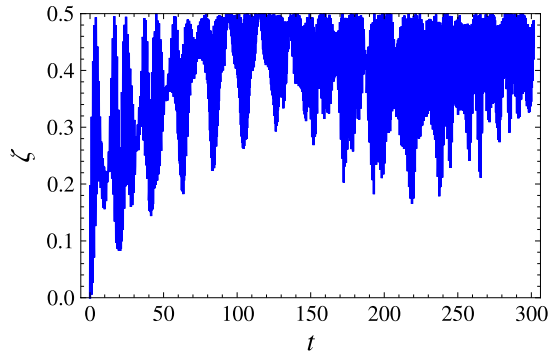


Fig. 6. Plot of the linear entropy ζ (as a function of t) of qubit1 initially in state $|e\rangle$ and the oscillator initially in a binomial state with $M = 11$ and $q = 0.95$. In this case qubit2 is coupled to the oscillator, with $\lambda_2 = 0.1$, $\lambda_1 = 1.0$, and $p = 0.5$.

Fano factor $F = 0.1$, relatively close to a number state, and it is coupled to a maximally mixed “environment”, $p = 0.5$. We note a significant departure from a regular oscillatory behaviour. Thus, the joint influence of the oscillator energy fluctuations and the weakly coupled “environment” (qubit2) is quite apparent even for longer times, as shown in Fig. 6.

5. Conclusion

In this work, I have considered the time evolution of the linear entropy of a two-level system (qubit1) coupled to an oscillator when the latter is interacting with a noisy two-level system (qubit2). This simple model has analytical solution within the RWA, which allowed a non-perturbative treatment and the accurate investigation of the dynamics of the system at all times. Nevertheless, despite being an integrable system, a relatively complex dynamics may arise; this is of course related to peculiar features of the dynamics, the linear entropy containing terms with frequencies that are in general incommensurate, which results in a quasi periodic dynamics. Besides, uncertainties in the “environment” (quantified by its degree of mixedness) have a significant influence on the qubit1 dynamics. In particular, the evolution of the linear entropy is highly sensitive to variations in the parameter p , or the initial mixedness of qubit2. Therefore, despite of the fact that the system under consideration is non-chaotic, it might show a (quasi periodic) very irregular evolution; the reversibility represented by the perfect periodicity of the Rabi oscillations will be strongly affected by a noisy but small “environment”. I should recall that such behaviour is different from the one in which qubit2 is not coupled ($\lambda_2 = 0.0$) and the oscillator is initially in a mixed state, as discussed in Section 2. There are a few similarities, but the presence of a third subsystem (qubit2) brings about, for instance, coherent exchange of excitations between the system and its “environment”, and a more complex dynamics will arise. I have also considered the effect of additional sources of noise, e.g., energy fluctuations in the initial preparation of the oscillator; even a small spread in energy together with the action of the “small environment” will significantly contribute to the loss of purity of qubit1.

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